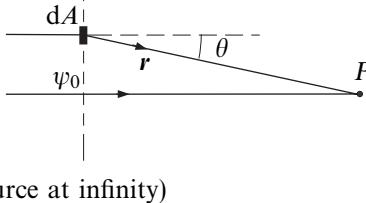
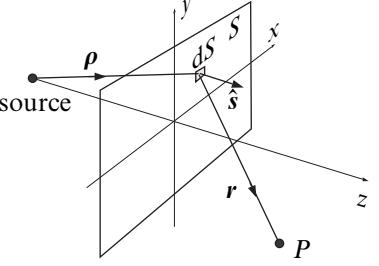


## 8.4 Fresnel diffraction

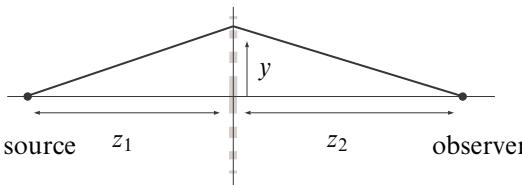
### Kirchhoff's diffraction formula<sup>a</sup>

 (source at infinity)		<p><math>\psi_P</math> complex amplitude at <math>P</math>  <math>\lambda</math> wavelength  <math>k</math> wavenumber (<math>= 2\pi/\lambda</math>)  <math>\psi_0</math> incident amplitude  <math>\theta</math> obliquity angle  <math>r</math> distance of <math>dA</math> from <math>P</math> (<math>\gg \lambda</math>)  <math>dA</math> area element on incident wavefront  <math>K</math> obliquity factor  <math>dS</math> element of closed surface  <math>\hat{s}</math> unit vector  <math>s</math> vector normal to <math>dS</math>  <math>r</math> vector from <math>P</math> to <math>dS</math>  <math>\rho</math> vector from source to <math>dS</math>  <math>E_0</math> amplitude (see footnote)</p>
Source at infinity $\psi_P = -\frac{i}{\lambda} \psi_0 \int K(\theta) \frac{e^{ikr}}{r} dA \quad (8.45)$ <p>where:</p> Obliquity factor (cardioid) $K(\theta) = \frac{1}{2}(1 + \cos\theta) \quad (8.46)$		
Source at finite distance <sup>b</sup> $\psi_P = -\frac{iE_0}{\lambda} \oint \frac{e^{ik(\rho+r)}}{2\rho r} [\cos(\hat{s} \cdot \hat{r}) - \cos(\hat{s} \cdot \hat{\rho})] dS \quad (8.47)$		

<sup>a</sup>Also known as the “Fresnel–Kirchhoff formula.” Diffraction by an obstacle coincident with the integration surface can be approximated by omitting that part of the surface from the integral.

<sup>b</sup>The source amplitude at  $\rho$  is  $\psi(\rho) = E_0 e^{ik\rho}/\rho$ . The integral is taken over a surface enclosing the point  $P$ .

### Fresnel zones

	<p><math>z</math> effective distance  <math>z_1</math> source-aperture distance  <math>z_2</math> aperture-observer distance  <math>n</math> half-period zone number  <math>\lambda</math> wavelength  <math>y_n</math> <math>n</math>th half-period zone radius  <math>z_m</math> distance of <math>m</math>th zero from aperture  <math>R</math> aperture radius</p>
Effective aperture distance <sup>a</sup> $\frac{1}{z} = \frac{1}{z_1} + \frac{1}{z_2} \quad (8.48)$	
Half-period zone radius $y_n = (n\lambda z)^{1/2} \quad (8.49)$	
Axial zeros (circular aperture) $z_m = \frac{R^2}{2m\lambda} \quad (8.50)$	

<sup>a</sup>I.e., the aperture–observer distance to be employed when the source is not at infinity.

## Cornu spiral

Fresnel integrals <sup>a</sup>	$C(w) = \int_0^w \cos \frac{\pi t^2}{2} dt \quad (8.51)$ $S(w) = \int_0^w \sin \frac{\pi t^2}{2} dt \quad (8.52)$
Cornu spiral	$CS(w) = C(w) + iS(w) \quad (8.53)$ $CS(\pm\infty) = \pm \frac{1}{2}(1+i) \quad (8.54)$
Edge diffraction	$\psi_P = \frac{\psi_0}{2^{1/2}} [CS(w) + \frac{1}{2}(1+i)] \quad (8.55)$ <p>where <math>w = y \left( \frac{2}{\lambda z} \right)^{1/2}</math> <math>(8.56)</math></p>
Diffraction from a long slit <sup>b</sup>	$\psi_P = \frac{\psi_0}{2^{1/2}} [CS(w_2) - CS(w_1)] \quad (8.57)$ <p>where <math>w_i = y_i \left( \frac{2}{\lambda z} \right)^{1/2}</math> <math>(8.58)</math></p>
Diffraction from a rectangular aperture	$\psi_P = \frac{\psi_0}{2} [CS(v_2) - CS(v_1)] \times [CS(w_2) - CS(w_1)] \quad (8.59)$ <p>where <math>v_i = x_i \left( \frac{2}{\lambda z} \right)^{1/2}</math> <math>(8.60)</math></p> <p>and <math>w_i = y_i \left( \frac{2}{\lambda z} \right)^{1/2}</math> <math>(8.61)</math></p> <p>and <math>w_i = y_i \left( \frac{2}{\lambda z} \right)^{1/2}</math> <math>(8.62)</math></p>
	<p>C Fresnel cosine integral S Fresnel sine integral</p> <p>CS Cornu spiral <math>v, w</math> length along spiral</p> <p><math>\psi_P</math> complex amplitude at <math>P</math> <math>\psi_0</math> unobstructed amplitude <math>\lambda</math> wavelength <math>z</math> distance of <math>P</math> from aperture plane [see (8.48)] <math>y</math> position of edge</p> <p>coherent plane waves</p> <p><math>x_i</math> positions of slit sides <math>y_i</math> positions of slit top/bottom</p>

<sup>a</sup>See also Equation (2.393) on page 45.

<sup>b</sup>Slit long in x.