

## 4.7 High energy and nuclear physics

### Nuclear decay

Nuclear decay law	$N(t) = N(0)e^{-\lambda t}$	(4.163)	$N(t)$ number of nuclei remaining after time $t$ $t$ time $\lambda$ decay constant
Half-life and mean life	$T_{1/2} = \frac{\ln 2}{\lambda}$	(4.164)	$T_{1/2}$ half-life
	$\langle T \rangle = 1/\lambda$	(4.165)	$\langle T \rangle$ mean lifetime
Successive decays $1 \rightarrow 2 \rightarrow 3$ (species 3 stable)			
	$N_1(t) = N_1(0)e^{-\lambda_1 t}$	(4.166)	
	$N_2(t) = N_2(0)e^{-\lambda_2 t} + \frac{N_1(0)\lambda_1(e^{-\lambda_1 t} - e^{-\lambda_2 t})}{\lambda_2 - \lambda_1}$	(4.167)	$N_1$ population of species 1 $N_2$ population of species 2 $N_3$ population of species 3
	$N_3(t) = N_3(0) + N_2(0)(1 - e^{-\lambda_2 t}) + N_1(0) \left( 1 + \frac{\lambda_1 e^{-\lambda_2 t} - \lambda_2 e^{-\lambda_1 t}}{\lambda_2 - \lambda_1} \right)$	(4.168)	$\lambda_1$ decay constant $1 \rightarrow 2$ $\lambda_2$ decay constant $2 \rightarrow 3$
Geiger's law <sup>a</sup>	$v^3 = a(R - x)$	(4.169)	$v$ velocity of $\alpha$ particle $x$ distance from source $a$ constant
Geiger–Nuttall rule	$\log \lambda = b + c \log R$	(4.170)	$R$ range $b, c$ constants for each series $\alpha, \beta$ , and $\gamma$

<sup>a</sup>For  $\alpha$  particles in air (empirical).

### Nuclear binding energy

Liquid drop model <sup>a</sup>	$B = a_v A - a_s A^{2/3} - a_c \frac{Z^2}{A^{1/3}} - a_a \frac{(N-Z)^2}{A} + \delta(A)$	(4.171)	$N$ number of neutrons $A$ mass number ( $= N+Z$ ) $B$ semi-empirical binding energy $Z$ number of protons $a_v$ volume term ( $\sim 15.8$ MeV) $a_s$ surface term ( $\sim 18.0$ MeV) $a_c$ Coulomb term ( $\sim 0.72$ MeV) $a_a$ asymmetry term ( $\sim 23.5$ MeV) $a_p$ pairing term ( $\sim 33.5$ MeV)
	$\delta(A) \simeq \begin{cases} +a_p A^{-3/4} & Z, N \text{ both even} \\ -a_p A^{-3/4} & Z, N \text{ both odd} \\ 0 & \text{otherwise} \end{cases}$	(4.172)	$M(Z, A)$ atomic mass $M_H$ mass of hydrogen atom $m_n$ neutron mass
Semi-empirical mass formula	$M(Z, A) = Z M_H + N m_n - B$	(4.173)	

<sup>a</sup>Coefficient values are empirical and approximate.

## Nuclear collisions

Breit–Wigner formula <sup>a</sup>	$\sigma(E) = \frac{\pi}{k^2} g \frac{\Gamma_{ab}\Gamma_c}{(E-E_0)^2 + \Gamma^2/4}$ (4.174)	$\sigma(E)$ cross-section for $a+b \rightarrow c$
	$g = \frac{2J+1}{(2s_a+1)(2s_b+1)}$ (4.175)	$k$ incoming wavenumber $g$ spin factor $E$ total energy (PE + KE) $E_0$ resonant energy $\Gamma$ width of resonant state $R$ $\Gamma_{ab}$ partial width into $a+b$ $\Gamma_c$ partial width into $c$ $\tau$ resonance lifetime $J$ total angular momentum quantum number of $R$ $s_{a,b}$ spins of $a$ and $b$ $\frac{d\sigma}{d\Omega}$ differential collision cross-section $\mu$ reduced mass $K =  \mathbf{k}_{\text{in}} - \mathbf{k}_{\text{out}} $ (see footnote) $r$ radial distance $V(r)$ potential energy of interaction
Total width	$\Gamma = \Gamma_{ab} + \Gamma_c$ (4.176)	
Resonance lifetime	$\tau = \frac{\hbar}{\Gamma}$ (4.177)	
Born scattering formula <sup>b</sup>	$\frac{d\sigma}{d\Omega} = \left  \frac{2\mu}{\hbar^2} \int_0^\infty \frac{\sin Kr}{Kr} V(r) r^2 dr \right ^2$ (4.178)	
Mott scattering formula <sup>c</sup>	$\frac{d\sigma}{d\Omega} = \left( \frac{\alpha}{4E} \right)^2 \left[ \csc^4 \frac{\chi}{2} + \sec^4 \frac{\chi}{2} + \frac{A \cos \left( \frac{\alpha}{\hbar v} \ln \tan^2 \frac{\chi}{2} \right)}{\sin^2 \frac{\chi}{2} \cos \frac{\chi}{2}} \right]$ (4.179)	$\hbar$ (Planck constant)/ $2\pi$ $\alpha/r$ scattering potential energy $\chi$ scattering angle $v$ closing velocity $A = 2$ for spin-zero particles, $= -1$ for spin-half particles
	$\frac{d\sigma}{d\Omega} \simeq \left( \frac{\alpha}{2E} \right)^2 \frac{4 - 3 \sin^2 \chi}{\sin^4 \chi}$ ( $A = -1$ , $\alpha \ll v\hbar$ ) (4.180)	

<sup>a</sup>For the reaction  $a+b \rightarrow R \rightarrow c$  in the centre of mass frame.

<sup>b</sup>For a central field. The Born approximation holds when the potential energy of scattering,  $V$ , is much less than the total kinetic energy.  $K$  is the magnitude of the change in the particle's wavevector due to scattering.

<sup>c</sup>For identical particles undergoing Coulomb scattering in the centre of mass frame. Nonidentical particles obey the Rutherford scattering formula (page 72).

## Relativistic wave equations<sup>a</sup>

Klein–Gordon equation (massive, spin zero particles)	$(\nabla^2 - m^2)\psi = \frac{\partial^2 \psi}{\partial t^2}$ (4.181)	$\psi$ wavefunction $m$ particle mass $t$ time
Weyl equations (massless, spin 1/2 particles)	$\frac{\partial \psi}{\partial t} = \pm \left( \boldsymbol{\sigma}_x \frac{\partial \psi}{\partial x} + \boldsymbol{\sigma}_y \frac{\partial \psi}{\partial y} + \boldsymbol{\sigma}_z \frac{\partial \psi}{\partial z} \right)$ (4.182)	$\psi$ spinor wavefunction $\boldsymbol{\sigma}_i$ Pauli spin matrices (see page 26)
Dirac equation (massive, spin 1/2 particles)	$(i\gamma^\mu \partial_\mu - m)\psi = 0$ (4.183) where $\partial_\mu = \left( \frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$ (4.184) $(\gamma^0)^2 = \mathbf{1}_4$ ; $(\gamma^1)^2 = (\gamma^2)^2 = (\gamma^3)^2 = -\mathbf{1}_4$ (4.185)	$i$ $i^2 = -1$ $\gamma^\mu$ Dirac matrices: $\gamma^0 = \begin{pmatrix} \mathbf{1}_2 & 0 \\ 0 & -\mathbf{1}_2 \end{pmatrix}$ $\gamma^i = \begin{pmatrix} 0 & \boldsymbol{\sigma}_i \\ -\boldsymbol{\sigma}_i & 0 \end{pmatrix}$ $\mathbf{1}_n$ $n \times n$ unit matrix

<sup>a</sup>Written in natural units, with  $c = \hbar = 1$ .